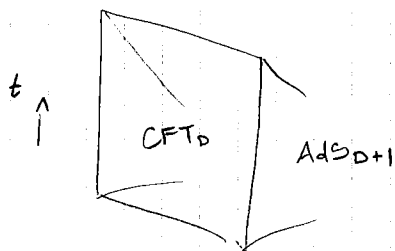
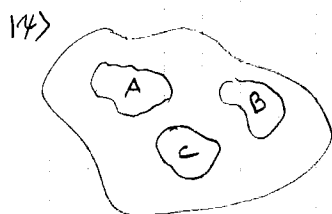


1) Intro: a holographic puzzle

- Recall, AdS/CFT conjecture: certain states of CFT_D dual to AdS_{D+1}



- Not all CFT states have a geometric dual, however.
- Natural Q: what fraction of states are holographic?
- One strategy: check holographic entanglement inequalities



let $\rho_A = \text{Tr}_A |\Psi\rangle\langle\Psi|$, etc.
 $S(A) = -\text{Tr} \rho_A \log \rho_A$ etc.

- All states obey, e.g. SSA: $S(AB) + S(BC) \geq S(ABC) + S(CB)$
- Not all states obey MMI $S(AB) + S(BC) + S(AC) \geq S(A) + S(B) + S(C) + S(ABC)$ but, holographic states do!

Method 1:

- randomly generate pure state of N qubits (in practice $N=5$)
- trace out k qubits
- check MMI for all perms of A, B, C st. $B = |ABC| \leq N-k$

Method 2:

- randomly generate entropy vector $(S_A, S_B, S_C, S_{AB}, S_{BC}, S_{AC}, S_{ABC})$
- check that it actually corresponds to a state
- check MMI

- ①: almost all states obey MMI
- ②: $\approx 1/2$ of states obey MMI

what gives?

A: Random States are not holographic!

2) Random States: Page's Theorem

Defⁿ A random state is the random variable $|\Psi(U)\rangle = U|1\rangle$ \leftarrow ref. state
 \leftarrow drawn from some prob. distrⁿ

Q: why useful?

- \rightarrow parameterize ignorance
- \rightarrow computationally useful/easier to avg. over many trials
- \rightarrow ??

Defⁿ Let $\mathcal{H}: \dim \mathcal{H} = n < \infty$. The normalized Haar measure on $U(n) \ni$ the unique measure μ s.t.

$$\begin{aligned} (1) \quad & \mu(SU) = \mu(US) = \mu(S) \quad \forall S \in U(n), U \in U(n) \\ (2) \quad & \mu(U(n)) = 1 \end{aligned}$$

$\rightarrow \mu$ defines a Haar-random state

Thm (Page) Suppose $\mathcal{H} = \mathcal{H}_A \otimes \mathcal{H}_B$, $e_A(U) := \text{Tr}_B |K(U)\rangle\langle K(U)|$. Then

$$\int d\mu(U) \|e_A(U) - \frac{I_A}{d_A}\|_1 \leq \sqrt{\frac{d_A}{d_B}} \quad \begin{aligned} d_A &= \dim \mathcal{H}_A \\ d_B &= \dim \mathcal{H}_B \end{aligned}$$

// Aside: Why $\|\cdot\|_1$? $\|\cdot\|_1: L(\mathcal{H}) \rightarrow \mathbb{R}$

// $T \mapsto \text{Tr} \sqrt{T^\dagger T}$

// If $\|e - \sigma\|_1 < \epsilon$, then $\|P(e - \sigma)\|_1 < \epsilon \quad \forall$ projectors $P \rightarrow$ measures distinctness

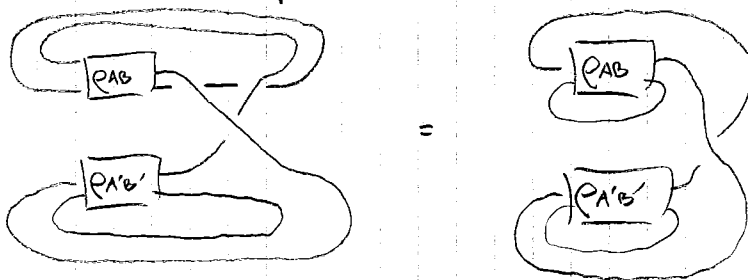
~~A slick proof in 4 steps:~~

① Let $e_A = \text{Tr}_B |K\rangle\langle K|_{AB}$, introduce a copy $A'B'$

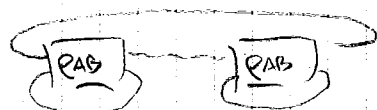
$$\text{Claim: } \text{Tr}_A e_A^2 = \text{Tr}_{ABA'B'} [(S_{AA'} \otimes I_{BB'}) |K\rangle\langle K|_{AB} \otimes |K\rangle\langle K|_{A'B'}]$$

\uparrow SWAP operator

pf. by picture: RHS



LHS



② Let $\langle \cdot \rangle \equiv$ Haar avg. over all pure $|K\rangle$

a result: $\langle |K\rangle\langle K|_A \otimes |K\rangle\langle K|_{A'} \rangle = C \Pi_{AA'}$
 a const \uparrow projector onto subspace that is symm under $A \leftrightarrow A'$

// could prove using properties of Haar μ , but should be fairly obvious
 // \rightarrow state being avg'd is symm
 // \rightarrow avg. cannot depend on any single state

q: what is C ?

$$\rightarrow \text{Tr}(C \Pi_{AA'}) = 1 \quad \text{so} \quad C = 1/\text{Tr} \Pi_{AA'}$$

$$\begin{aligned}
\text{Tr } \Pi_{AA'} &= \sum_{a, a'=1}^{d_A} \langle a | \langle a' | \Pi_{AA'} | a \rangle | a' \rangle \\
&= \sum_{a, a'} \langle a | \langle a' | \left[\frac{1}{2} (I_{AA'} + S_{AA'}) \right] | a \rangle | a' \rangle \\
&= \frac{1}{2} \left[\sum_{a, a'} \langle a | a \rangle \langle a' | a' \rangle + \sum_{a, a'} \langle a | a' \rangle \langle a' | a \rangle \right] \\
&= \frac{1}{2} \left[\sum_{a, a'} 1 + \sum_a 1 \right] \\
&= \frac{1}{2} d_A (d_A + 1)
\end{aligned}$$

③ Now for $|e\rangle_{AB}$,

$$\begin{aligned}
\langle \text{Tr } \rho_A^2 \rangle &= \langle \text{Tr}_{AB, A'B'} [S_{AA'} \otimes I_{BB'} |e\rangle\langle e|_{AB} \otimes |e\rangle\langle e|_{A'B'}] \rangle \\
&= \text{Tr}_{AB, A'B'} [S_{AA'} \otimes I_{BB'} \langle |e\rangle\langle e|_{AB} \otimes |e\rangle\langle e|_{A'B'} \rangle] \\
&= \text{Tr}_{AB, A'B'} [S_{AA'} I_{BB'} \cdot C \cdot \Pi_{AB, A'B'}] \\
&= \dots \\
&= \frac{d_A + d_B}{d_A d_B + 1}
\end{aligned}$$

$$\begin{aligned}
\textcircled{4} \quad \left\langle \left\| \rho_A - \frac{I_A}{d_A} \right\|_2^2 \right\rangle &= \left\langle \text{Tr} \left(\rho_A^2 - \frac{2}{d_A} \rho_A + \frac{I_A}{d_A^2} \right) \right\rangle \\
&= \langle \text{Tr } \rho_A^2 \rangle - \frac{2}{d_A} \langle \text{Tr } \rho_A \rangle + \frac{1}{d_A^2} \langle \text{Tr } I_A \rangle \\
&= \frac{d_A^2 - 1}{d_A (d_A d_B + 1)} = \frac{d_A^2}{d_A^2 d_B} = \frac{1}{d_B}
\end{aligned}$$

• For positive $f(x)$, $\langle \sqrt{f} \rangle = \sqrt{\langle f \rangle}$

$$\begin{aligned}
\text{// schematic: } \langle \sqrt{f} \rangle &= \sum p(x) \sqrt{f(x)} \\
&= \langle \vec{p}, \vec{\sqrt{f}} \rangle \\
&\leq \|\vec{p}\|_2 \|\vec{\sqrt{f}}\|_2 \quad (\text{Cauchy-Schwarz}) \\
&= \sqrt{\sum p(x)^2} \sqrt{\sum f(x)} \\
&\leq \sqrt{\sum f(x)} = \sqrt{\sum p(x) f(x)} = \sqrt{\langle f \rangle}
\end{aligned}$$

$$\therefore \left\langle \left\| \rho_A - \frac{I_A}{d_A} \right\|_2 \right\rangle \leq \frac{1}{\sqrt{d_B}}$$

• Cauchy-Schwarz again: $\|\cdot\|_1 \leq \sqrt{d} \|\cdot\|_2$

$$\Rightarrow \left\langle \left\| \rho_A - \frac{I_A}{d_A} \right\|_1 \right\rangle \leq \sqrt{\frac{d_A}{d_B}}$$

3) Back to the puzzle:

The resolution: Haar-typical states generically saturate balanced inequalities

$$S(AB) + S(BC) + S(AC) \quad S(A) + S(B) + S(C) + S(ABC)$$

$$\log 2 \cdot (\# \text{ qubits in } A + \# \text{ qubits in } B) \quad \text{etc...}$$

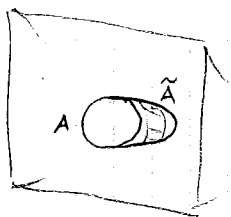
ex $N=5, k=1$
 $A: 1$ qubit
 $B: 1$ qubit
 $C: 2$ qubits

MHI: $2 + 2 + 2 > 1 + 1 + 2 + 1$

and importantly, Haar-typical states are not holographic!

PF Let $\mathcal{H}: \text{CFT}_D$, $|\psi\rangle$ a Haar-typical state, A : subregion w/ characteristic length ℓ (e.g. a $D-1$ ball)

$\rightarrow S(A)$ is extensive $S(\rho_A) \propto \left(\frac{\ell}{\epsilon}\right)^{D-1}$
 \uparrow UV cutoff



but, Ryu-Takayanagi: $S(\rho_A) = \frac{\text{area}(\tilde{A})}{4G}$

\rightarrow For suff. small A , \tilde{A} only probes asymp. AdS region

Exercise For $A \equiv D-1$ ball in CFT, area of spherical cap \tilde{A} in AdS is

$\text{area}(\tilde{A}) \lesssim \int_{\ell/\epsilon}^1 \frac{d\tilde{r}}{r^{D-1}}$ subextensive! #