## Demystifying Wheeler-de Witt

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## **1** Exercise: Minisuperspace

Consider the *minisuperspace* that consists of all closed homogeneous and isotropic FRW spacetimes in 1+3 dimensions that contain a homogeneous and isotropic real scalar field,  $\Phi$ . For fun, let's also suppose that there is a (non-negative) cosmological constant  $\Lambda$  that is fixed to some value, *i.e.*, we don't consider is as part of the minisuperspace. In a suitable coordinate system, the line element of the spacetime manifold reads

$$ds^2 = \frac{3}{\Lambda} \left( -dt^2 + a^2(t)d\Omega_3^2 \right),\tag{1}$$

where  $d\Omega_3^2$  is the round metric on the 3-sphere. In this setting, a 3-geometry and matter configuration is entirely characterized by only two real numbers: the value of the scale factor, b, and the value of the scalar field,  $\Phi_0$ . As such, the wavefunctional over geometries and matter configurations reduces to a wavefunction with two arguments:  $\Psi[\gamma_{ij}(\mathbf{x}), {\Phi(\mathbf{x})}] \equiv \Psi(b, \Phi_0)$ . In this exercise, we will deduce the form of the Wheeler-De Witt equation that this wavefunction obeys.

(a) In units where  $G = \hbar = c = 1$ , the Einstein-Hilbert action is

$$S_{EH} = -\frac{1}{16\pi} \int d^4x \sqrt{-g} \left(R - 2\Lambda\right) \tag{2}$$

and the action of a free scalar field is

$$S_{\Phi} = \frac{1}{2} \int d^4 x \sqrt{-g} \left( g^{\mu\nu} \Phi_{,\mu} \Phi_{,\nu} + m^2 \Phi^2 \right) \,. \tag{3}$$

Use the form of the metric above and the fact that  $\Phi(t)$  depends only on t to write down the total action  $S = S_{EH} + S_{\Phi}$ . You may find it convenient to express your answer in terms of the following quantities:

$$H^2 \equiv \frac{\Lambda}{3} \qquad \phi \equiv \sqrt{\frac{4\pi}{3}} \Phi \qquad \mu^2 \equiv \frac{3}{\Lambda} m^2 \tag{4}$$

(b) We can read off the Lagrangian  $\mathcal{L}$  from the total action  $S = \int dt \mathcal{L}$ . Take a and  $\phi$  to be the canonical coordinates and define their conjugate momenta via

$$p_a = \frac{\partial \mathcal{L}}{\partial \dot{a}} \qquad p_\phi = \frac{\partial \mathcal{L}}{\partial \dot{\phi}},\tag{5}$$

where a dot over a quantity denotes a derivative with respect to t. Perform a Legendre transformation and write down the Hamiltonian  $\mathcal{H}$  in terms of a,  $\phi$ ,  $p_a$ , and  $p_{\phi}$ . Note that  $\mathcal{H}$  has no explicit dependence on the coordinate time t; at any given time  $t_0$ ,  $\mathcal{H}$  depends only on the values of  $a(t_0) \equiv b$ ,  $\phi(t_0) \equiv \chi$ ,  $p_a(t_0) \equiv p_b$ , and  $p_{\phi}(t_0) \equiv p_{\chi}$ . As such, to emphasize that  $\mathcal{H}$  is purely a function of 3-geometry and the matter configuration, we can write  $\mathcal{H} = \mathcal{H}(b, \chi, p_b, p_{\chi})$ .

- (c) Note that there do not seem to be any secondary constraints for the momenta above, *i.e.*,  $\dot{p}_a, \dot{p}_\phi \neq 0$ . Why is this? Is it still correct to write  $\mathcal{H} = 0$ ?
- (d) Make the identification

$$p_b \to -i\frac{\partial}{\partial b} \qquad p_\chi \to -i\frac{\partial}{\partial \chi}$$
 (6)

and write down the explicit form of the Wheeler-de Witt equation  $\mathcal{H}\Psi(b,\chi) = 0$ .

(e) Express the Wheeler-De Witt operator  $\mathcal{H}(b,\chi,\partial_b,\partial_\chi)$  in the following form:

$$\mathcal{H} = \frac{1}{2} G^{AB} \partial_A \partial_B + \mathcal{V} \tag{7}$$

 $G_{AB}$  and  $\mathcal{V}$  are the Wheeler-De Witt metric and potential for minisuperspace in the coordinates  $x^A = (b, \chi)$ . What is the signature of the Wheeler-De Witt metric?